

Introduction to Medical Imaging

Cone-Beam CT

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Introduction

Available cone-beam reconstruction methods:

- exact
- approximate

Our discussion:

- exact (now)
- approximate (next)

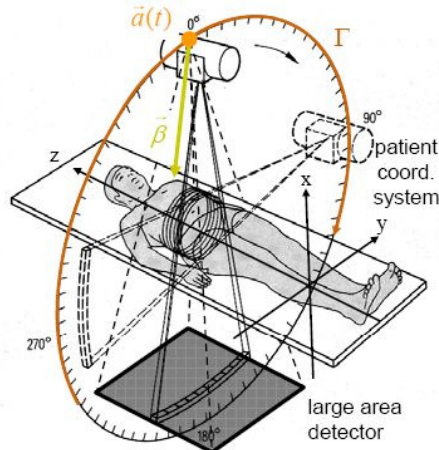
The Radon transform and its inverse are important mechanisms to understand cone-beam CT

Cone-Beam Transform

$$D\mu(\vec{a}(t), \vec{\beta}) = \int_0^\infty \mu(\vec{a}(t) + s\vec{\beta}) ds, \quad (\vec{a}, \vec{\beta}) \in \Gamma \times S^2$$

$\vec{a}(t)$ is the source position along trajectory Γ
 $\vec{\beta}$ the unit vector pointing along a particular x-ray beam

The cone-beam transform reflects the data acquisition process of measuring line integrals of the attenuation coefficient μ .



2D Radon Transform

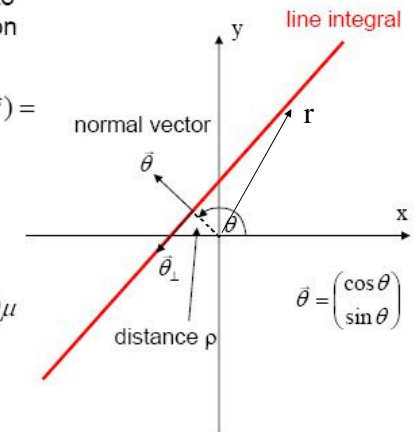
The analytical approach of reconstruction by projections has to be done in the context of the Radon transform \Re

$$\Re\mu(\rho, \vec{\theta}) = \int_{-\infty}^{+\infty} dl \mu(\rho \cdot \vec{\theta} + l \cdot \vec{\theta}_\perp) = \int_{-\infty}^{+\infty} dl \mu(\rho \cdot \vec{\theta} + l \cdot \vec{\theta}_\perp)$$

Thus in the 2D case the Radon transform $\Re\mu$ is identical to the measured cone beam transform $D\mu$

$$D\mu(\vec{a}, \vec{\theta}_\perp) \Big|_{\vec{a} \cdot \vec{\theta} = \rho} = \Re\mu(\rho, \vec{\theta})$$

with projection angle θ .

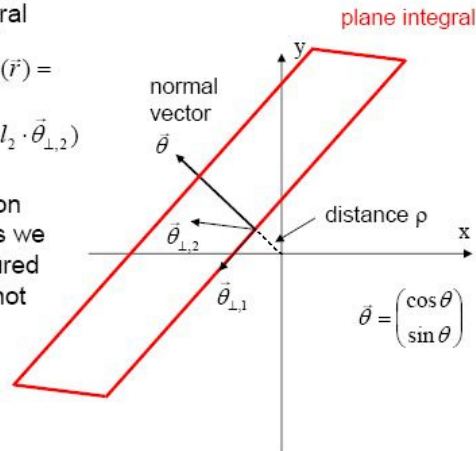


3D Radon Transform

In three dimensions the Radon transform \Re is a plane integral

$$\Re\mu(\rho, \vec{\theta}) = \int d^3r \delta(\vec{r} \cdot \vec{\theta} - \rho) \cdot \mu(\vec{r}) = \int_{-\infty}^{+\infty} dl_1 \int_{-\infty}^{+\infty} dl_2 \mu(\rho \cdot \vec{\theta} + l_1 \cdot \vec{\theta}_{\perp,1} + l_2 \cdot \vec{\theta}_{\perp,2})$$

which is a severe complication compared to the 2D case. As we will see the link to the measured cone beam transform $D\mu$ is not trivial.

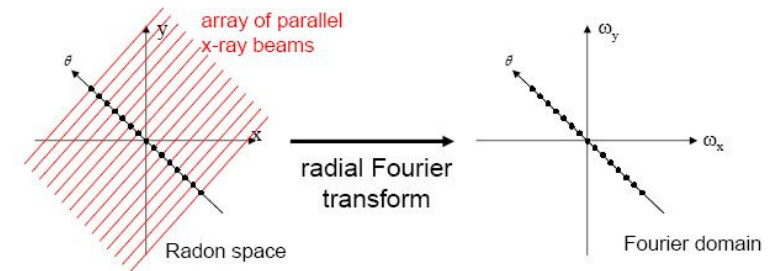


from: Dr. Günter Lauritsch, Siemens

Fourier-Slice Theorem in 2D

$$F_{\rho} \Re\mu(\rho, \vec{\theta}) = (F_2\mu)(\omega_{\rho} \cdot \vec{\theta})$$

The radial 1D Fourier transform F_{ρ} of the Radon transform $\Re\mu$ along $\vec{\theta}$ is equal to the 2D Fourier transform F_2 of the object μ along $\vec{\theta}$ perpendicular to the direction of the projection.

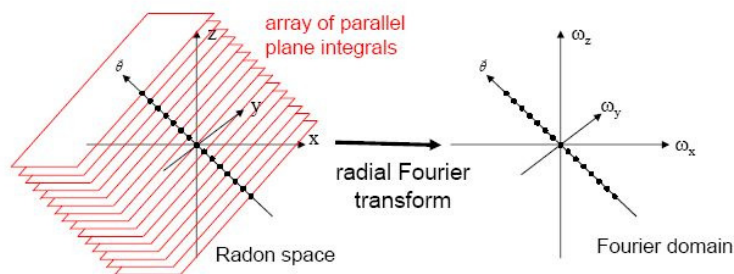


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Fourier-Slice Theorem in 3D

$$F_{\rho} \Re\mu(\rho, \vec{\theta}) = (F_3\mu)(\omega_{\rho} \cdot \vec{\theta})$$

The radial 1D Fourier transform F_{ρ} of the Radon transform $\Re\mu$ along $\vec{\theta}$ is equal to the 3D Fourier transform F_3 of the object μ along $\vec{\theta}$ perpendicular to the direction of the projection.



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Exact Reconstruction in 2D and 3D

In 2D:

- use 2D inversion formula: the filtered backprojection procedure
- we have seen a spatial technique, only performing filtering in the frequency domain (in a polar grid)
- but may also interpolate the polar grid in the frequency domain and invert the resulting cartesian lattice
- employ linogram techniques for the latter (see later)

In 3D:

- use 3D inversion formula: not nearly as straightforward than 2D inversion
- full frequency-space methods also exist
- more details next (on all)

Exact Inversion Formula

The basic 3D inversion filtered backprojection formula, due to Natterer (1986):

$$f(\mathbf{x}) = \frac{-1}{8\pi^2} \int_{S^2} \frac{\partial^2}{\partial \rho^2} \Re f(|\rho|\theta) d\theta.$$

- θ is the angle, a unit vector on a unit sphere
- \mathbf{x} , ρ are object and Radon space coordinates, resp.: $|\rho| = \mathbf{x} \cdot \theta$
- involves a 2nd derivative of the 3D Radon transform
- the second derivative operator can be treated as a convolution kernel

Some manipulations can reduce the second derivative to a first derivative, along with convolution operators

$$f(\mathbf{x}) = \frac{1}{2} \int_{S^2} \frac{-1}{4\pi^2} \frac{\partial^2}{\partial \rho^2} \Re f(|\rho|\theta) d\theta = \frac{1}{2} \int_{S^2} \frac{-1}{2\pi^2 \rho^2} * \frac{\partial}{\partial \rho} \left[\frac{1}{2\pi^2 \rho} * \Re f(|\rho|\theta) \right] d\theta$$

- many different variants have been proposed
 - for example: Kudo/Saito (1990), Smith (1985)

Grangeat's Algorithm

Phase 1:

- from cone-beam data to derivatives of Radon data

Phase 2:

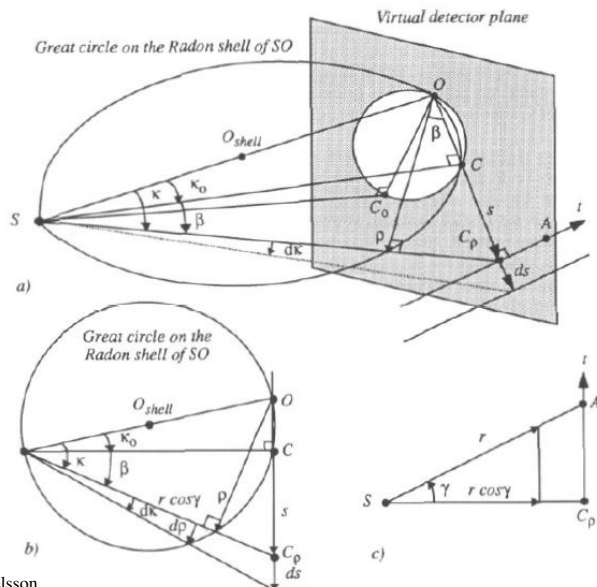
- from derivatives of Radon data to reconstructed 3D object

There are many ways to achieve Phase 2

- direct, $O(N^5)$
- a two-step procedure, $O(N^4)$ [Marr et al, 1981]
- a Fourier method, $O(N^3 \log N)$, [Axelsson/Danielsson, 1994]
- a divide-and-conquer strategy, $O(N^3 \log N)$ [Basu/Bresler, 2002]
- we shall discuss the first three here

But first let us see how Radon data are generated from cone-beam data

Transforming Cone-Beam to Radon Data



from Axelsson/Danielsson

Transforming Cone-Beam to Radon Data

$$\begin{aligned} \frac{d}{d\rho} [\Re f(\rho)] &= \int_{-\pi/2}^{\pi/2} \int_0^\infty \frac{d}{d\rho} f(\rho, r, \gamma) r dr d\gamma = \frac{d}{d\kappa} \int_{-\pi/2}^{\pi/2} X f(\rho, \gamma) \frac{1}{\cos \gamma} d\gamma \\ &= \frac{SC}{\cos^2 \beta} \frac{d}{ds} \int_{-\infty}^\infty \frac{1}{SA} X f(\rho, t) dt. \end{aligned}$$

Strategy:

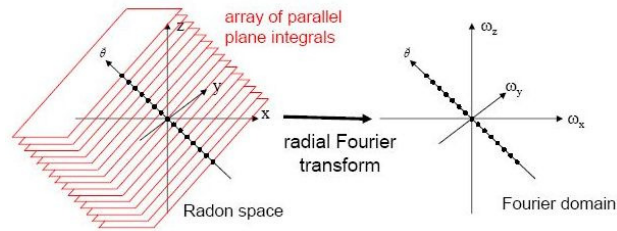
- weigh detector data with a factor $1/SA$
- integrate along all intersections (lines) between the detector plane and the required Radon planes
 - there are N^2 such lines (N lines and N rotations)
- take the derivative in the s -direction (in the detector plane perpendicular to t)
- weight the 2D data set resulting from a single source position by the factor $SC / \cos^2 \beta$

The order of these operations can be switched since they are all linear (Grangeat swapped the order of operation 2 and 3)

Radon Data to Object: Direct Method

There are $O(N^3)$ data points in Radon (derivative) space

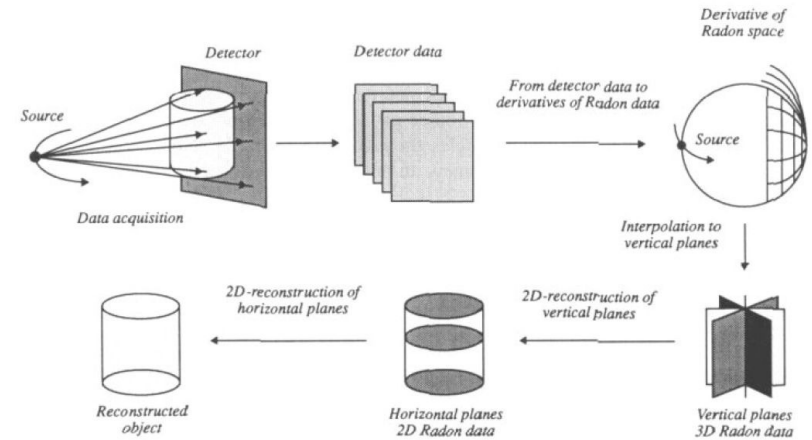
Each is due to a plane integral



The direct method simply inserts the plane data into the object space, one by one

- this is basically the expansion of a point into a plane, defined by (θ, ρ)
- this gives rise to an $O(N^5)$ algorithm

Radon Data to Object: Two-Step Method



from Axelsson/Danielsson

Radon Data to Object: Two-Step Method

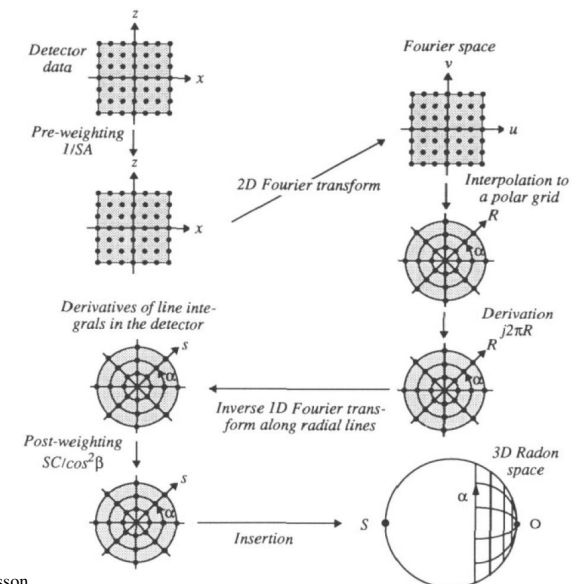
Each vertical plane holds all Radon points due to plane integrals of perpendicularly intersecting planes

- filtered backprojection reduces the plane integrals to line integrals, confined to horizontal planes

The horizontal planes are then reconstructed with another filtered backprojection

Each such operation is $O(N^3)$ and there are $O(N)$ of them, resulting in a complexity of $O(N^4)$

Radon Data to Object: Fourier Space Approach



from Axelsson/Danielsson

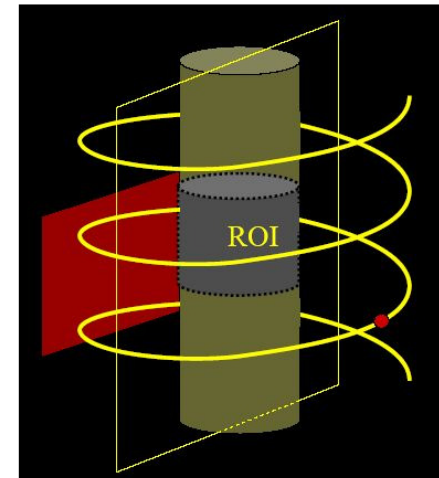
Radon Data to Object: Fourier Space Approach

Takes advantage of the $O(N \log N)$ complexity of the FFT at various steps

It also uses linograms [Edholm/Herman, 1987] to reduce 2D interpolation to 1D interpolation

The complexity is then $O(N^3 \log N)$

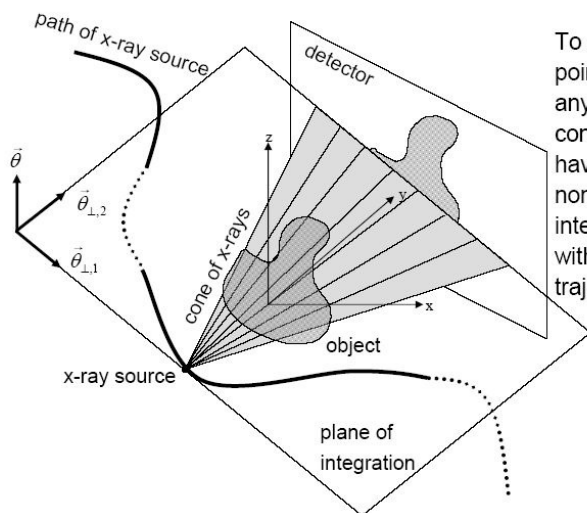
Long Object Problem



- Reconstruction of an ROI should be feasible from projection data restricted to the ROI and some surrounding.
- The basic 3D Radon inversion formula does not fulfill this request.

from: Dr. Günter Lauritsch, Siemens

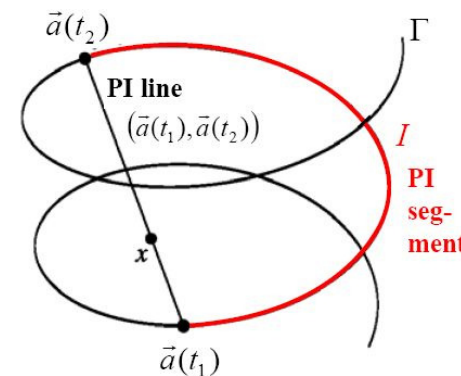
Tuy's Sufficiency Condition



To reconstruct a point x of the object any plane containing x must have at least one non tangential intersection point with the source trajectory.

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Concept of PI-Lines

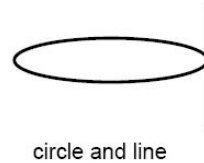
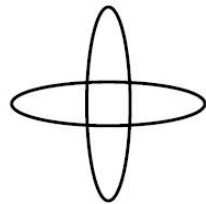
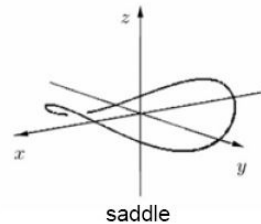
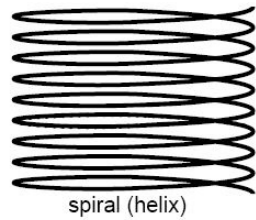


For a point x on a PI line any plane containing x has at least one intersection point with the PI segment associated with the PI line.

The PI segment is that portion of the source trajectory needed for reconstructing the point x .

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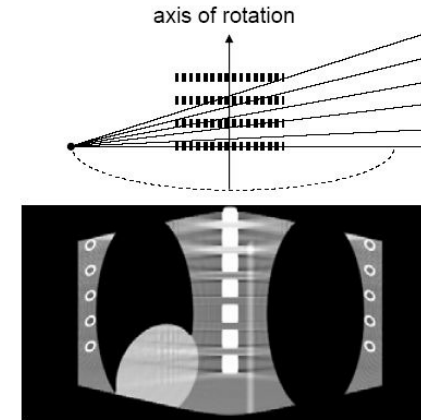
Examples of Complete Trajectories



from: Dr. Günter Lauritsch, Siemens

Circular Source Path

A prominent example of an incomplete trajectory

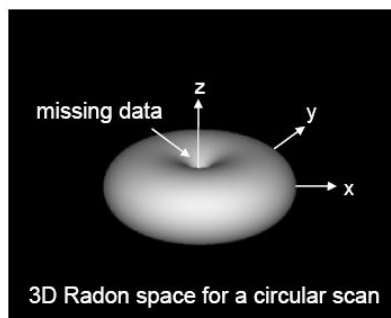


- Due to incomplete data sampling cone artifacts show up at sharp z-edges of objects with high contrast.
- Almost horizontal rays (or integration planes) are missing to distinguish between the members of the object stack.

Thorax simulation study.
Coronal slice. $C=0$, $W=200$

from: Dr. Günter Lauritsch, Siemens

3D Radon Data Acquired by a Circular Trajectory



By a circular source trajectory a donut shaped region is acquired in 3D Radon space. At the z-axis a cone-like region is missing.

from: Dr. Günter Lauritsch, Siemens

Challenges in Cone-Beam Reconstruction

The naive application of the 3D Radon inversion formula is prohibitive due to

- long object problem
- enormous computational expense

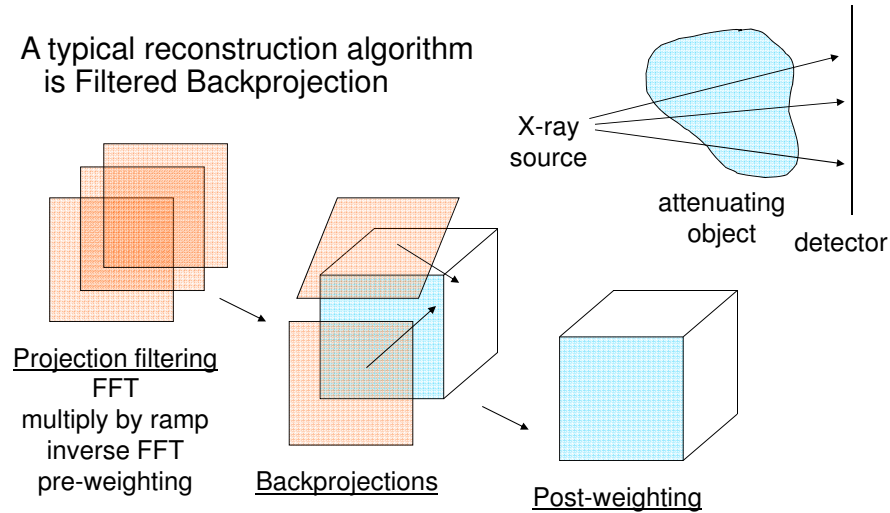
Simplifications have to be found to end up in an efficient and numerically stable reconstruction algorithm preferably in a shift-invariant 1D-filtered backprojection algorithm

Utilization of redundant data is obscure. Ideally redundancy in collected Radon planes has to be considered. However, this approach is suboptimal because:

- it is quite complicated
- underestimates the redundancy of data
- typically in cone beam, the data are highly redundant in approximation

Transmission CT

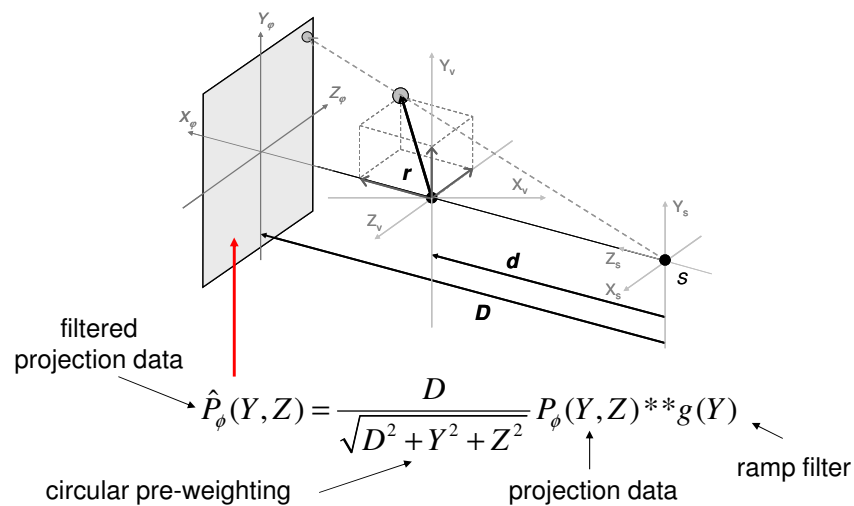
A typical reconstruction algorithm is Filtered Backprojection



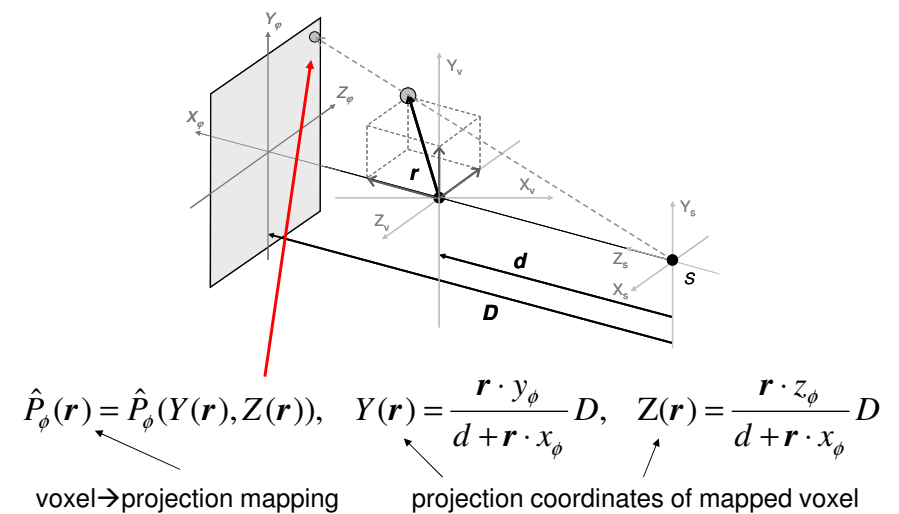
Popular Approximation

Feldkamp-Davis-Kress (FDK) Cone-beam reconstruction

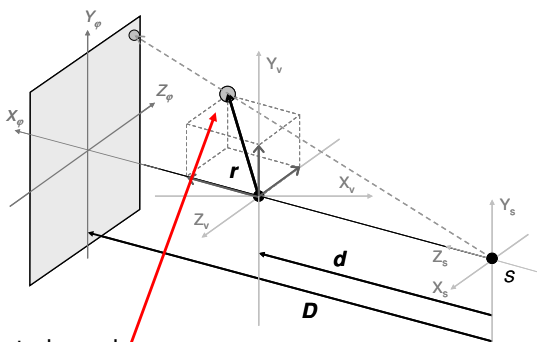
FDK: Filtering



FDK: Backprojection



FDK: Accumulation, Depth-Weighting



reconstructed voxel

$$f(\mathbf{r}) = \frac{1}{4\pi^2} \int_0^{2\pi} \frac{d^2}{(d + \mathbf{r} \cdot \mathbf{x}_\phi)^2} \hat{P}_\phi(\mathbf{r}) d\phi$$

accumulation for all projections

depth-weighting